

1 A mechanism-denial study on the Madden-Julian
2 Oscillation with reduced interference from mean
3 state changes

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4 Mechanism-denial experiments using Superparameterized Community At-
5 mosphere Model are conducted to investigate the importance of extratrop-
6 ical and circumnavigating waves, wind-evaporation feedback and radiative-
7 convective feedback to the Madden-Julian Oscillation (MJO). A common is-
8 sue with mechanism-denial studies is the interference from mean state changes
9 when processes are turned off in the model. Here, time-invariant forcing and
10 nudging on effective timescales longer than the intraseasonal timescale are
11 implemented to maintain the mean state. The MJO activity remains largely
12 unchanged with suppressed extratropical and circumnavigating waves when
13 the mean state is maintained to be close to that of the control run, suggest-
14 ing that excitation of MJO by extratropical and circumnavigating waves is
15 not necessary for the existence of MJO in this model. It is also shown that
16 the wind-evaporation feedback slows down eastward propagation of the MJO,
17 and the radiative-convective feedback amplifies the MJO.

1. Introduction

18 The Madden-Julian oscillation [MJO; Madden and Julian, 1971] is the dominant mode
19 of intraseasonal variability in the tropics, and has been extensively studied in the past
20 few decades [see, e.g., review by Zhang, 2005]. The characteristics of the MJO are well
21 documented. The MJO features a spatial scale of zonal wavenumbers 1-3 and eastward
22 propagation primarily in the Indian Ocean and western Pacific Ocean at around 5 m s^{-1} .
23 Below we introduce the processes considered important to MJO initiation and maintenance
24 in previous studies and examined here, including (i) surface heat fluxes and radiative
25 fluxes, i.e., sources of the column-integrated moist static energy (hereafter column MSE)
26 [e.g., Hu and Randall, 1994; Maloney et al., 2010; Andersen and Kuang, 2012], and (ii)
27 influences of extratropical waves [e.g., Liebmann and Hartmann, 1984; Lau and Phillips,
28 1986] and circumnavigating waves [e.g., Knutson and Weickmann, 1987; Matthews, 2008].

29 The MJO has been hypothesized as a moisture mode in the sense that column moisture
30 dominates the column MSE variations, and the growth and maintenance of the MJO is
31 often interpreted in the context of processes recharging the column MSE while the MJO's
32 propagation is sometimes viewed as the result of horizontal moisture advection which has
33 been argued to lead column MSE [e.g., Fuchs and Raymond, 2002; Maloney et al., 2010;
34 Sobel and Maloney, 2012; Sobel et al., 2014; Pritchard and Bretherton, 2014]. Wind-
35 evaporation feedback and radiative-convective feedback are examples of such processes
36 and are considered important to the MJO [e.g., Bony and Emanuel, 2005; Sobel et al.,
37 2010]. With details different from the originally proposed linear theory [Emanuel, 1987;
38 Neelin et al., 1987], surface heat flux anomalies, led by surface wind anomalies, are argued

39 to be positively correlated with the column MSE, and therefore strengthen the MJO [e.g.,
40 Maloney, 2009; Sobel et al., 2010; Kiranmayi and Maloney, 2011]. The radiative heating
41 anomaly is found to be in phase with MJO precipitation and column MSE [Lin and Mapes,
42 2004; Ma and Kuang, 2011], and plays a significant role in maintaining the column MSE
43 anomaly associated with the MJO [e.g., Andersen and Kuang, 2012; Sobel et al., 2014].

44 The recharge-discharge mechanisms for the MJO are internal to the tropics. However,
45 coherence between the tropical convection and extratropical circulation on the intrasea-
46 sonal timescale has been noticed since the 1980s [e.g., Weickmann, 1983; Lau and Phillips,
47 1986], and Hsu [1996] suggested that the intraseasonal oscillation is a global phenomenon.
48 Despite controversy over the statistical significance of the dependence between the trop-
49 ical and extratropical intraseasonal signals [Ghil and Mo, 1991], many studies tried to
50 better understand the interactions between the oscillations in the tropics and those in
51 the mid-latitudes. One direction of the interaction is through the propagation of Rossby
52 waves generated by tropical heating, which leads to global responses [Jin and Hoskins,
53 1995; Matthews et al., 2004]. In the other direction, it has been observed that extratrop-
54 ical waves can propagate into the tropics in the regions with westerlies and lead to MJO
55 initiation [e.g., Hsu et al., 1990]. Hoskins and Yang [2000] further showed that westerly
56 winds are not necessary for extratropical influence on the tropical atmosphere because
57 extratropical forcing can project onto the equatorial modes. More recent results from
58 numerical integrations indicate that proper extratropical forcing leads to MJO initiation
59 [Lin et al., 2007; Ray et al., 2009; Ray and Zhang, 2010]. In addition to extratropical
60 waves, it has also been proposed that eastward circum-equatorial propagation of a Kelvin

61 wave caused by a preceding MJO event [e.g., Knutson and Weickmann, 1987; Matthews,
62 2008; Maloney, 2015].

63 Mechanism-denial experiments have been conducted in recent years to examine the im-
64 portance of particular processes to the MJO [Kim et al., 2011; Ray and Li, 2013; Pritchard
65 and Bretherton, 2014]. However, in these mechanism-denial studies, MJO response to the
66 suppression of a particular process without being mediated by mean state changes is not
67 distinguished from the response to the departure of mean state away from the control sim-
68 ulations, which leads to ambiguity in the interpretation. To examine the direct influences
69 (i.e., without being mediated by mean state changes) of particular processes on the MJO,
70 we have designed measures, largely inspired by Hall [2000], to minimize changes to the
71 mean state when processes are disabled in the model. The methodology will be introduced
72 in Section 2, including the model configuration, and experimental setup. In Section 3, we
73 will first examine how the climatology and MJO activity respond in the experiments with
74 suppressed extratropical waves, and use these experiments as an example to illustrate the
75 importance of maintaining the mean state in mechanism-denial studies. The results from
76 the other experiments are presented in the same section. A brief summary and further
77 discussions follow in Section 4.

2. Methodology

2.1. Model

78 We used the Super-Parameterized Community Atmosphere Model [SPCAM; Khairout-
79 dinov and Randall, 2001] Version 3.5, in which the conventional cloud parameterizations
80 in the Community Atmosphere Model are replaced by a two-dimensional cloud system-

81 resolving model [Khairoutdinov and Randall, 2003]. SPCAM is chosen because it is
82 known to produce robust MJO signals without tuning convection parameterization [e.g.,
83 Khairoutdinov et al., 2005; Andersen and Kuang, 2012; Benedict et al., 2014; Pritchard
84 and Bretherton, 2014]. The simulations are conducted using semi-Lagrangian advection
85 at T42 resolution (around 280 km) with 30 vertical levels on the Community Atmosphere
86 Model level. On each grid point of the Community Atmosphere Model component, the
87 embedded cloud-resolving array has 32 grid points that spans 128 km in the north-south
88 direction. All the integrations are forced with perpetual February, the peak season of
89 MJO activities [Zhang and Dong, 2004], sea surface temperature (Hadley Center Opti-
90 mally Interpolated Sea Surface Temperature, averaged from 1980-2000). The simulations
91 are integrated for 13 years and the first 3 years are discarded.

2.2. Experimental setup

92 Table 1 provides a brief description of the experiments. Inspired by Hall [2000], time-
93 invariant forcing is applied to maintain the mean state. For each experiment, an ensemble
94 of 28 simulations are restarted from different points of the CTL run, and integrated for
95 20 days. The first ensemble member restarts from the point where the CTL run has been
96 integrated for six years, with each of the of following ensembles lagging the previous one
97 by three months. Time-invariant forcing is computed using the tendency with which the
98 ensemble mean drifts away from the climatology of the CTL run, and then applied to
99 the next set of ensemble simulations. After four such iterations, the forcing can mostly
100 compensate for the effects of the processes turned off. We also nudged the prognostic
101 variables in the tropics back to the control run climatology on effective timescales longer

102 than the intraseasonal timescale to make sure the mean state is well maintained in a
103 long-term integration (see Supporting Information for details).

104 To constrain the impact of the nudging described in the Supporting Information, NDG
105 is integrated with the prognostic variables nudged back to the mean state of CTL. In
106 CHN and CHNn, the intraseasonal disturbance can grow and propagate within a wide
107 channel in the tropics, as the prognostic variables are relaxed towards the control run
108 climatology north of 30°N and south of 40°S . The relaxation rate starts from zero at
109 30°N and 40°S , and increases linearly poleward till 42°N and 52°S , respectively, so that
110 the waves emanating from the tropics are not reflected back. Poleward of 42°N and 52°S ,
111 the relaxation time is fixed at 9 hours. These latitudes are chosen so that the damping
112 reaches sufficiently low latitudes to adequately suppress the extratropical waves, but does
113 not extend too far into the tropics and suppress the patterns (e.g., Rossy gyres) associated
114 with the MJO. The procedure to maintain the climatology is applied to all the experiments
115 discussed in this paper, except CHNn (and CTL), so the climatology of the CHNn run
116 can be substantially different from that of the CTL run. Comparison of CHNn and CHN
117 runs highlights the interference due to changes to the climatology. To test the sensitivity
118 to the latitudes beyond which the extratropical waves are suppressed, the equatorward
119 boundaries of the relaxation zones are shifted from 30°N and 40°S to 10°N and 30°S in
120 CHN2, leaving a narrower channel compared to CHN. In the experiment BOX, we further
121 relax the prognostic variables toward the control run climatology all over the globe except
122 within the box over the tropical Indian and Pacific Oceans (40°S - 30°N , 0°E - 270°E), so
123 that both extratropical waves and circumnavigating waves are suppressed. In addition,

124 radiative-convective (wind-evaporation) feedback is turned off by prescribing the radiative
125 heating (surface fluxes) using the control run climatology in the RAD (FLUX) experiment.
126 Again, the climatologies of all experiments except CHNn (and CTL) are maintained to
127 be that of the CTL run.

3. Results

128 The CTL run simulates the climatology reasonably well compared to observations. The
129 precipitation maxima are located in the western Indian Ocean, the Maritime Continent
130 and the western Pacific Ocean between 20°S and the Equator, as it is forced by perpetual
131 February sea surface temperature (Figure 1A). The time-space power spectrum of the
132 Outgoing Longwave Radiation (OLR) is calculated with successive 90-day segments of
133 data overlapping by 60 days in the fashion of Wheeler and Kiladis [1999]. The power
134 spectrum of precipitation is calculated in the same way and shown in the Supporting
135 Information, confirming that OLR is a good proxy for deep convection. The peak at
136 small wavenumbers and low frequencies in the spectrum shows a robust MJO signal in
137 the CTL run (Figure 2A). The MJO bandpass filtered (30-90 days, wavenumbers 1-3)
138 variance of OLR (Figure 1A) and precipitation (Figure S2A) indicate that the peak of MJO
139 activity spans the Indian Ocean and the western Pacific Ocean between 20°S-8°S, a few
140 degrees poleward compared to observations during boreal winter [e.g., Zhang, 2005]. When
141 the prognostic variables are nudged on effective timescales longer than the intraseasonal
142 timescale, the mean precipitation remains similar to that of the CTL run. While the
143 tropical transients of the lowest resolved frequency (1/90 cpd) are damped (Figure 2B),
144 the spatial distribution of MJO activity is also largely unchanged (Figure 1B). Despite

145 that the power spectrum is dominated by a flat peak at wavenumbers 1-4, the frequency-
146 wavenumber diagram confirms that the simulation produces a strong MJO signal when
147 the nudging is implemented (Figure 2B). Regressed MJO precipitation and 850 hPa wind
148 anomalies also show realistic life cycle of MJO (see Supporting Information for details).

149 We now turn to the responses of the mean states and MJO activity in the experiments
150 with suppressed extratropical waves, and these experiments will serve as an example to
151 highlight the importance of maintaining the mean states. In CTL, there are large values
152 of eddy momentum flux convergence in the midlatitudes (Figure 3A), as Rossby waves
153 originating from these latitudes propagate into the tropics and the polar regions. When
154 the extratropical waves are suppressed, the eddy momentum flux convergence disappears
155 in the Northern Hemisphere, while there is still weak eddy momentum flux convergence
156 on the equator flank of the damping zone (the equatorward boundary is marked by the
157 dashed line) in the Southern Hemisphere (Figure 3B). When the extratropical waves are
158 damped starting from lower latitudes, there is a single peak of eddy momentum flux
159 convergence in the upper atmosphere lying over the tropics (Figure 3C), and it confirms
160 that the extratropical waves are adequately suppressed. In fact, the responses of the MJO
161 are similar in both the wide channel (CHN) and the narrow channel (CHN2) experiments,
162 so only results from CHN will be shown hereafter.

163 The CTL run produces a strong cross-equatorial winter cell (Figure 4A), and the mag-
164 nitude of the Hadley Cells matches reanalysis data well [Schneider, 2006]. When the
165 extratropical waves are suppressed, the Hadley Cells and Ferrel Cells weaken (Figure 4B)
166 if the mean state is not maintained (the CHNn run). In contrast, there is hardly any

167 changes in the circulation over the tropics in CHN, and it suggests that the method im-
168 plemented in the present study is indeed able to minimize the departure of the mean
169 states away from the CTL run. The response of the climatological precipitation shares a
170 similar picture. As the Hadley Cells weaken in CHNn, the precipitation shifts poleward.
171 The precipitation decreases over the Inter Tropical Convergence Zone, especially over the
172 Equatorial Indian Ocean, and increases over the South Pacific Convergence Zone (Fig-
173 ure 1C), while the precipitation in CHN remains similar to that of the CTL run. The
174 averaged precipitation over the domain with high MJO activity (between 25°S-10°N and
175 30°E-120°W) is 6.0 mm day^{-1} in CTL, which is maintained to be the same in CHN, and
176 reduced by 15% in CHNn. The pattern correlation between the precipitation in CTL and
177 CHNn is 0.916, compared to 0.995 between that in CTL and CHN. Although the MJO
178 signal seems to change little in the wavenumber-frequency spectrum (Figure 2C), the OLR
179 variance of MJO decreases significantly in the Indian Ocean in CHNn (Figure 1C). As long
180 as the climatology does not change, the spatial distribution of the MJO activity remains
181 broadly unchanged even when the extratropical waves are suppressed (Figure 1D). The
182 power spectrum also confirms that there is a strong signal of MJO without extratropical
183 influence (Figure 2D). The regressed MJO precipitation and low-level winds anomalies are
184 concentrated over western Pacific in CHNn, while the MJO life cycle in CHN compares
185 well with that in NDG and CTL (Figure S4-5).

186 The mean states are well maintained in BOX, FLX and RAD, in which the averaged
187 precipitation is 6.0, 6.1 and 5.9 mm day^{-1} , and the pattern correlation of the precipitation
188 with that in CTL is 0.994, 0.993 and 0.994, respectively. The influences of the circum-

189 navigating waves, wind-evaporation feedback and radiative-convective feedback can thus
190 be assessed with reduced interference from mean state changes. Suppressing extratropical
191 waves and circumnavigating waves barely changes the spatial distribution of MJO activity
192 (Figure 1E). As the disturbances are suppressed over the Atlantic sector, the OLR power
193 spectrum is dominated by a peak at zonal wavenumbers 2-3 (Figure 2E). The essential
194 processes of the MJO are internal to the Equatorial Indian Ocean and the Pacific Ocean
195 in the experiment BOX, as MJO initiation and maintenance are limited within these re-
196 gions in the BOX experiment. The activity of the eastward propagating intraseasonal
197 disturbances with planetary scales weakens significantly in FLX and RAD (Figure 1FG),
198 which implies that the wind-evaporation feedback and radiative-convective feedback are
199 important to the MJO. In FLX, the spectral peak shifts to higher frequency compared
200 to the other experiments (Figure 2F), which indicates that the wind-evaporation feed-
201 back slows down MJO propagation in SPCAM. The MJO signal largely disappears in
202 the power spectrum of the OLR when the radiative feedback is disabled (Figure 2G), in
203 agreement with Kim et al. [2011]. On the other hand, the Kelvin wave becomes stronger
204 without interactive radiative heating in RAD, consistent with the results from Andersen
205 and Kuang [2012].

4. Summary and discussions

206 MJO mechanism-denial experiments have been conducted in a few recent studies [e.g.,
207 Ray and Li, 2013], in which, particular processes are disabled in numerical models, and the
208 importance of these processes to the MJO and evaluated based on how MJO responds in
209 the simulations. The response of MJO in these studies can be decomposed into two parts:

210 (1) response to mean state changes and (2) direct response to the suppression of certain
211 process without being mediated by mean state changes. Using SPCAM, the present
212 study seeks to examine the direct influences of the extratropical and circumnavigating
213 waves, wind-evaporation feedback and radiative-convective feedback on the MJO with
214 reduced interference from mean state changes. Using time-invariant forcing, combined
215 with nudging the prognostic variables back to the control run climatology on effective
216 timescales longer than the intraseasonal timescale, we were able to minimize the changes
217 to the climatology from CTL when processes are eliminated or suppressed in the model.
218 Without maintaining the climatology, the strength of the general circulation and tropical
219 precipitation is reduced in response to suppressed extratropical waves as seen in Ray and
220 Li [2013]. In particular, precipitation and MJO activity decrease significantly over the
221 Indian Ocean. It is worth noting here the response of the mean state and MJO activity
222 to the suppression of extratropical waves in CHNn is somewhat different from that in the
223 simulations in Ray and Li [2013], which may be attributed to, besides model difference,
224 the fact that the simulations are forced with perpetual February sea surface temperature
225 in the present study, while the simulations in Ray and Li [2013] are integrated with
226 annual cycle. The MJO activity remains generally the same as in the CTL run when
227 the mean state change is minimized, leading to the conclusion that excitation of MJO by
228 extratropical eddies are not essential for the existence of the MJO. This contrast highlights
229 the importance of reducing the interference from mean state changes in mechanism-denial
230 experiments. Results also suggest that the processes important to the MJO are internal to
231 the tropical Indian and Pacific Oceans in SPCAM, because the simulation produces strong

232 MJO signals when the transients can only grow and propagate within these regions. The
233 wind-evaporation feedback is found to slow down MJO propagation in the model. When
234 the radiative-convective feedback is turned off, the MJO signal weakens significantly.

235 As the results show that excitation of MJO by extratropical and circumnavigating waves
236 are not necessary for MJO's existence, it leads to the open question that how the present
237 results can be reconciled with previous studies, in which the influences of extratropical
238 and circumnavigating waves are argued to be important to MJO initiation. It is possible
239 that the extratropical waves and circumnavigating waves determine the timing for some
240 MJO initiation events, but the collective effects of the waves over a long period have little
241 influence on MJO climatology. Besides, the circumnavigating waves might play a more
242 significant role in MJO amplitude in idealized models such as that used by Maloney [2015].
243 The wind-evaporation feedback is shown to slow down the MJO, which is consistent with
244 reanalysis data and numerical studies [Kiranmayi and Maloney, 2011; Maloney, 2009].
245 In an idealized simulation of the MJO [Andersen and Kuang, 2012], the surface flux is
246 found to project positively onto the tendency of the column MSE, and the difference is
247 attributed to the mean surface easterly winds in the idealized simulation, compared to
248 the weak to westerly winds in the observations and realistic simulations. One caveat with
249 the present study is the simulations are forced with fixed sea surface temperature, while
250 it has been argued that including air-sea coupling in SPCAM can produce more realistic
251 phasing between surface flux and precipitation and improve the simulated MJO [Benedict
252 et al., 2011]. The radiative-convective feedback amplifies the MJO not only because, in
253 the active phase of the MJO, the anomalous column-integrated radiative heating serves as

254 a direct source of column MSE, but also due to the bottom-heavy profile of the radiative
255 heating anomaly [Ma and Kuang, 2011], which generates a bottom-heavy vertical velocity
256 profile and can further import MSE into the column. In contrast, such bottom-heavy
257 radiative heating anomaly in the convectively active phase weakens the Kelvin wave, as
258 the stratiform instability for convectively coupled waves counts on top-heavy convective
259 heating [Mapes, 2000; Kuang, 2008]. Also, a radiative-convective instability involving the
260 interaction among radiative heating, water vapor, clouds and large-scale circulation has
261 been proposed recently [Emanuel et al., 2014; Wing and Emanuel, 2014], and observational
262 and numerical studies are needed to relate this instability to realistic MJO initiations.

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Table 1. Brief description of the experiments. See text for details.

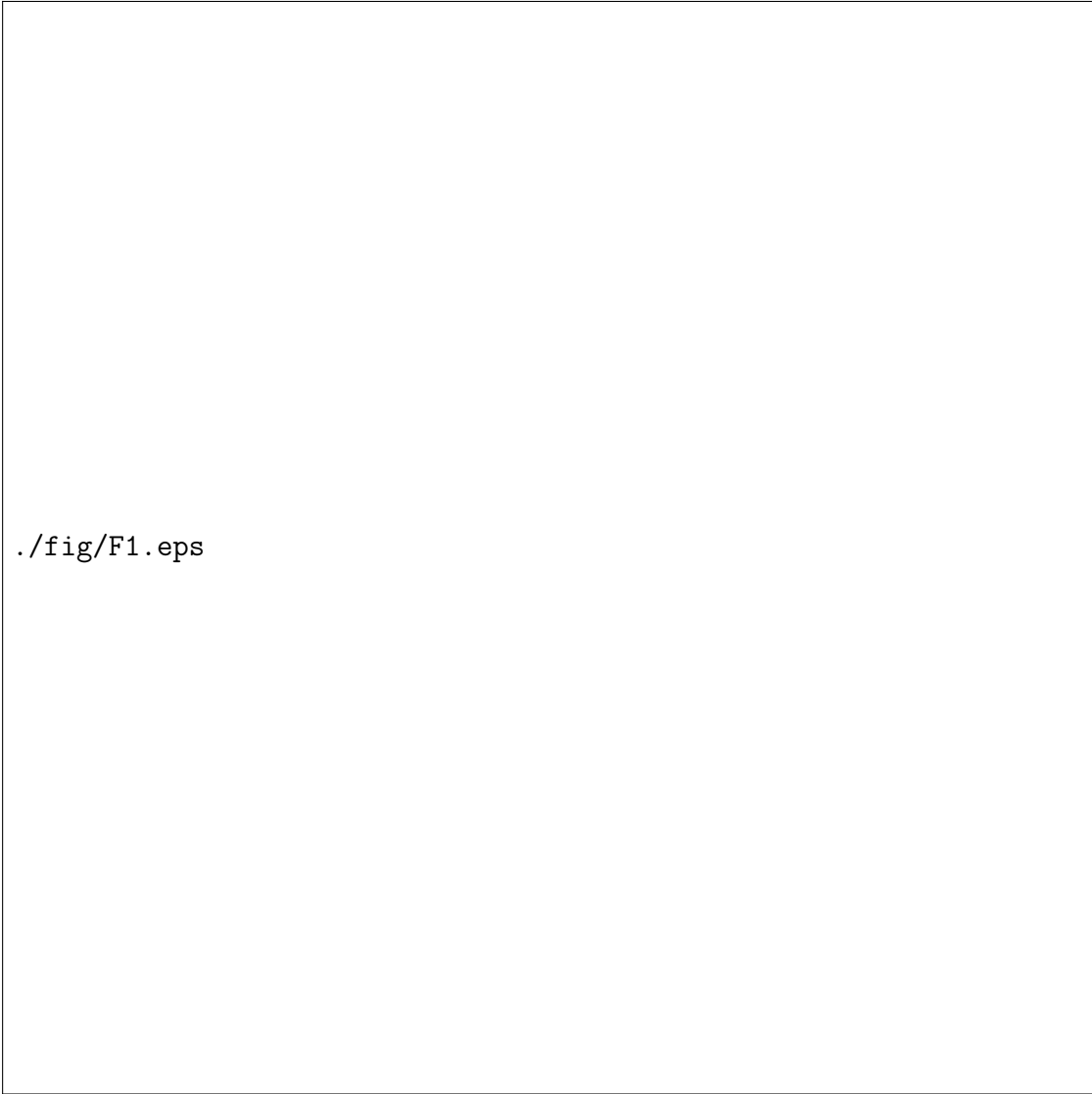
	process disabled	time-invariant forcing	nudging
CTL	n/a	×	×
NDG	n/a	×	✓
CHNn	extratropical waves outside of 30°N and 40°S	×	×
CHN	extratropical waves outside of 30°N and 40°S	✓	✓
CHN2	extratropical waves outside of 10°N and 30°S	✓	✓
BOX	extratropical wave outside of 30°N and 40°S and circumnavigating waves	✓	✓
FLX	wind-evaporation feedback	✓	✓
RAD	radiative-convective feedback	✓	✓

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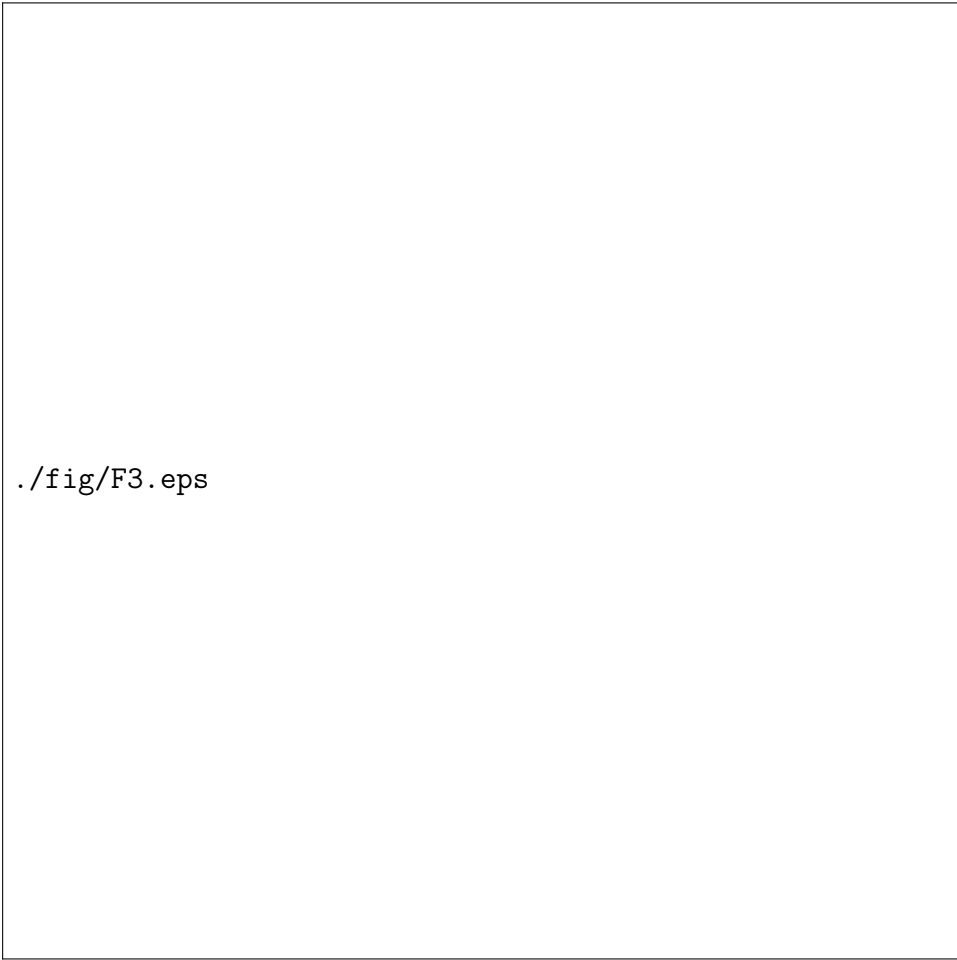


./fig/F1.eps

Figure 1. Mean precipitation (shaded) and OLR variance of the MJO (black contours) from the (A) CTL, (B) NDG, (C) CHN_n, (D) CHN, (E) BOX, (F) FLX and (G) RAD. The silver contours denote the coast lines.



Figure 2. Time-space power spectra of the OLR averaged between 20°S - 0° of (A) CTL, (B) NDG, (C) CHN_n, (D) CHN, (E) BOX, (F) FLX and (G) RAD.



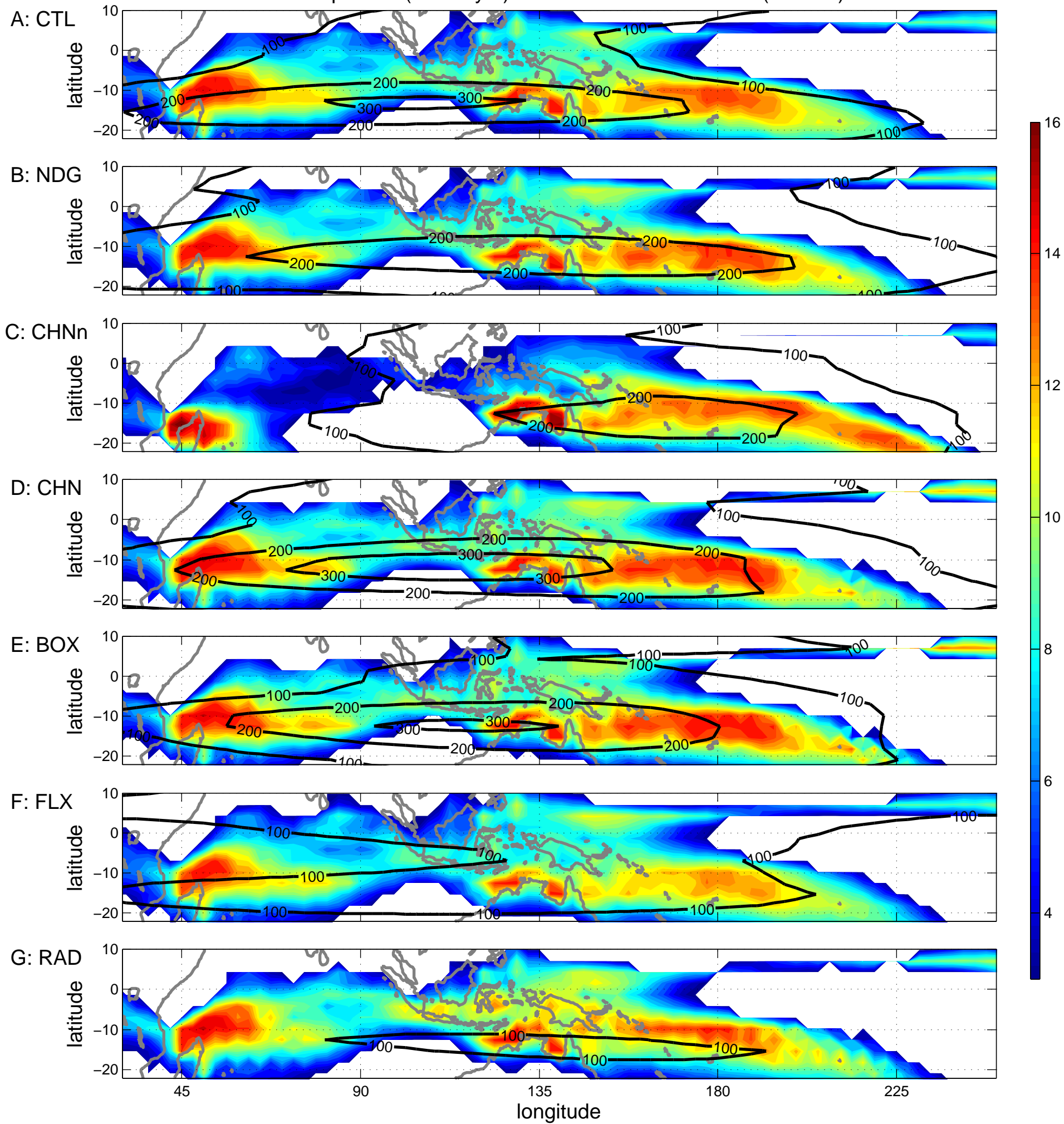
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Figure 3. Temporal and zonal mean eddy forcing in (A) CTL, (B) CHN and (C) CHN2. In (B) and (C), the dashed lines denote the latitudes starting from which the extratropical eddies are damped. The relaxation rate increases linearly poleward till the solid lines, and remains constant poleward from the solid lines.



Figure 4. Temporal and zonal mean circulation in (A) CTL, and anomalous circulation in (B) CHNn and (C) CHN. Negative streamfunction values (dashed contours) correspond to clockwise circulation, and positive streamfunction values (solid contours) correspond to counter-clockwise circulation.

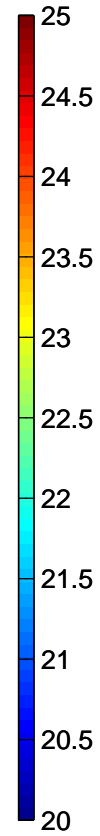
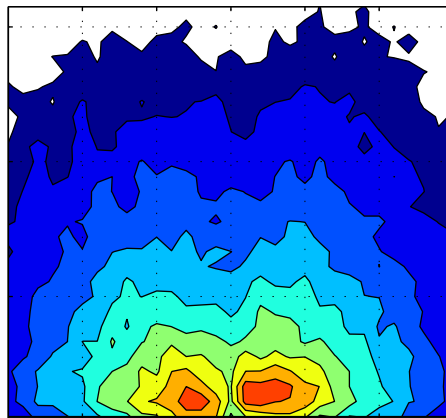
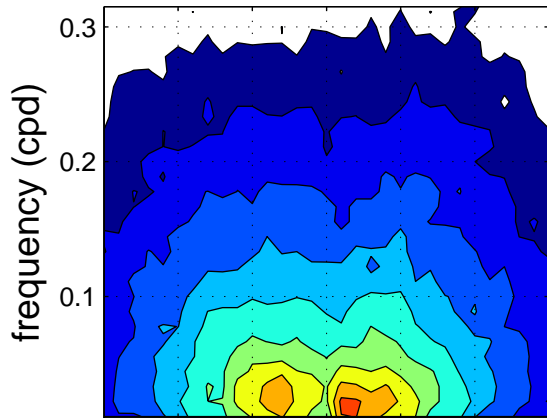
Precipitation (mm day^{-1}) and OLR variance of MJO ($\text{W}^2 \text{m}^{-4}$)



Power spectrum of OLR

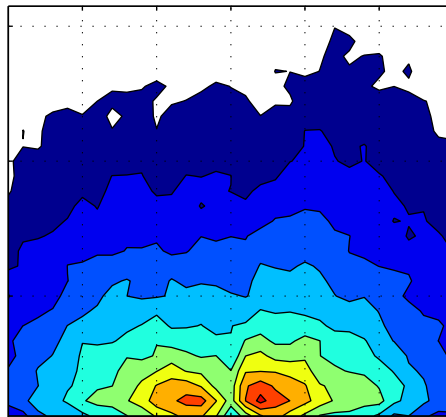
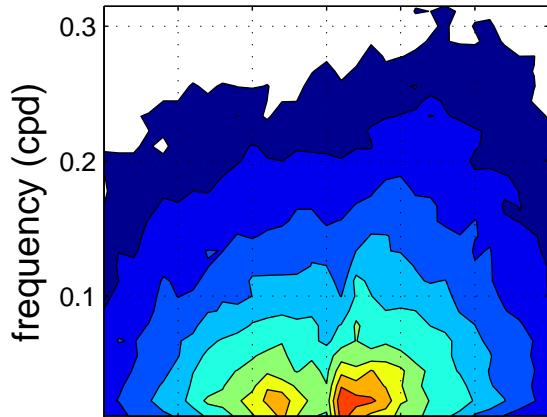
A: CTL

B: NDG



C: CHNn

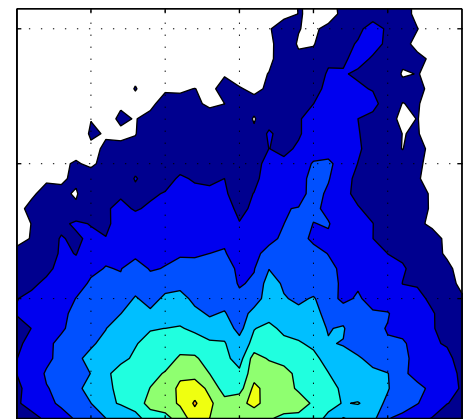
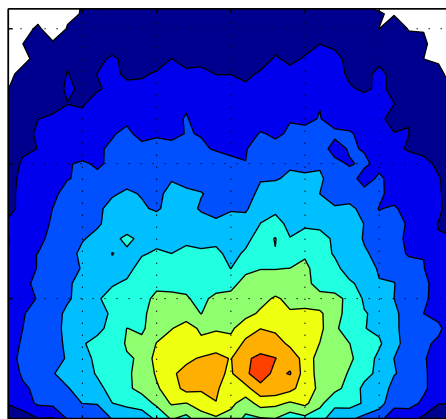
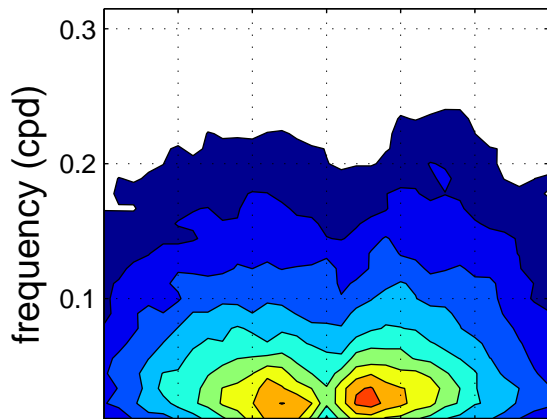
D: CHN



E: BOX

F: FLX

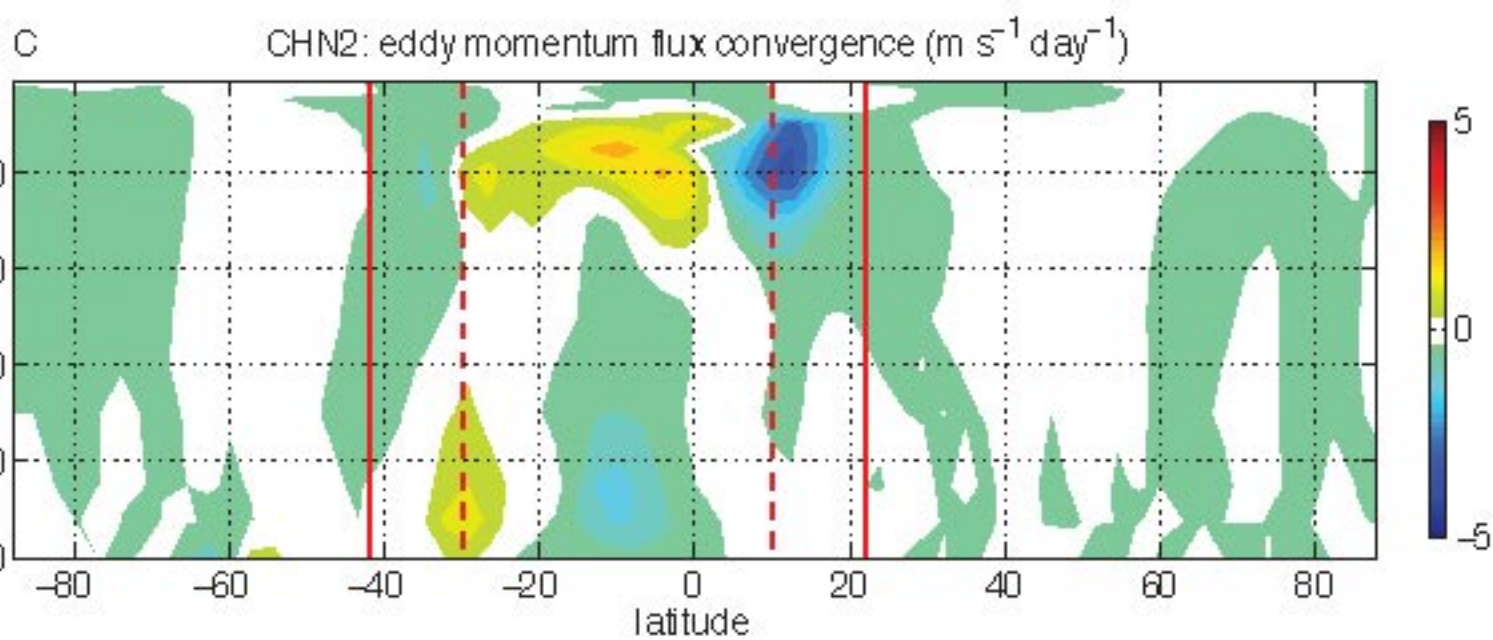
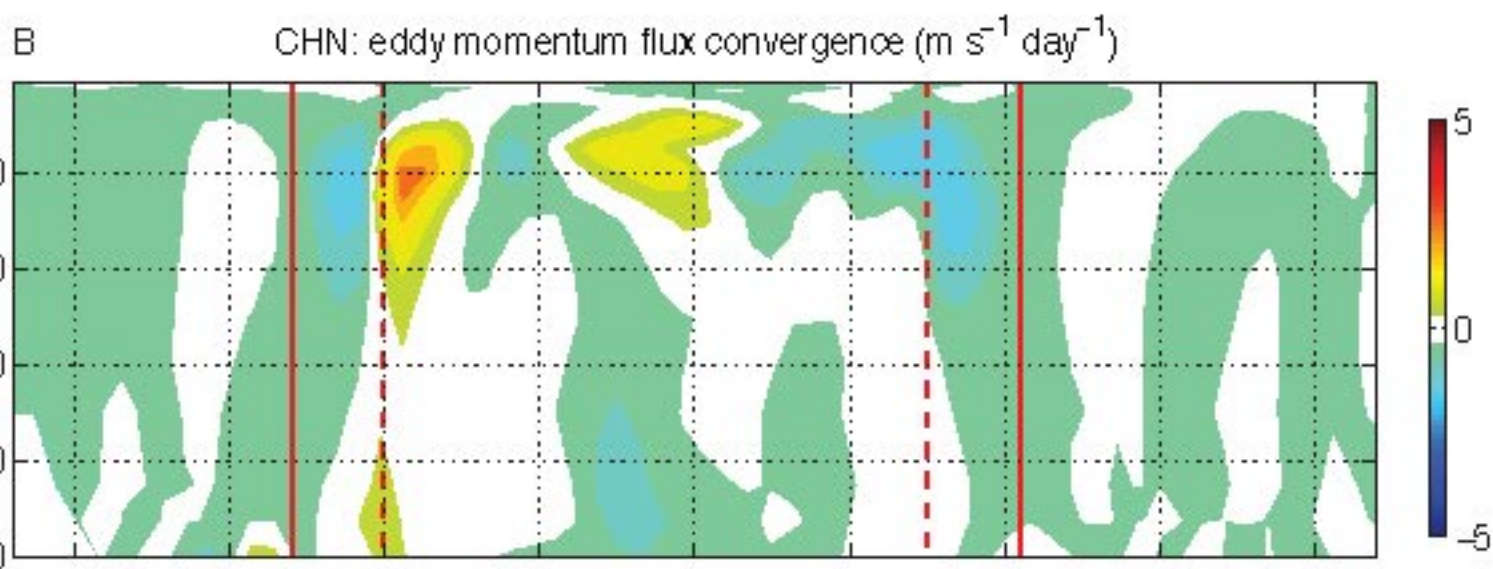
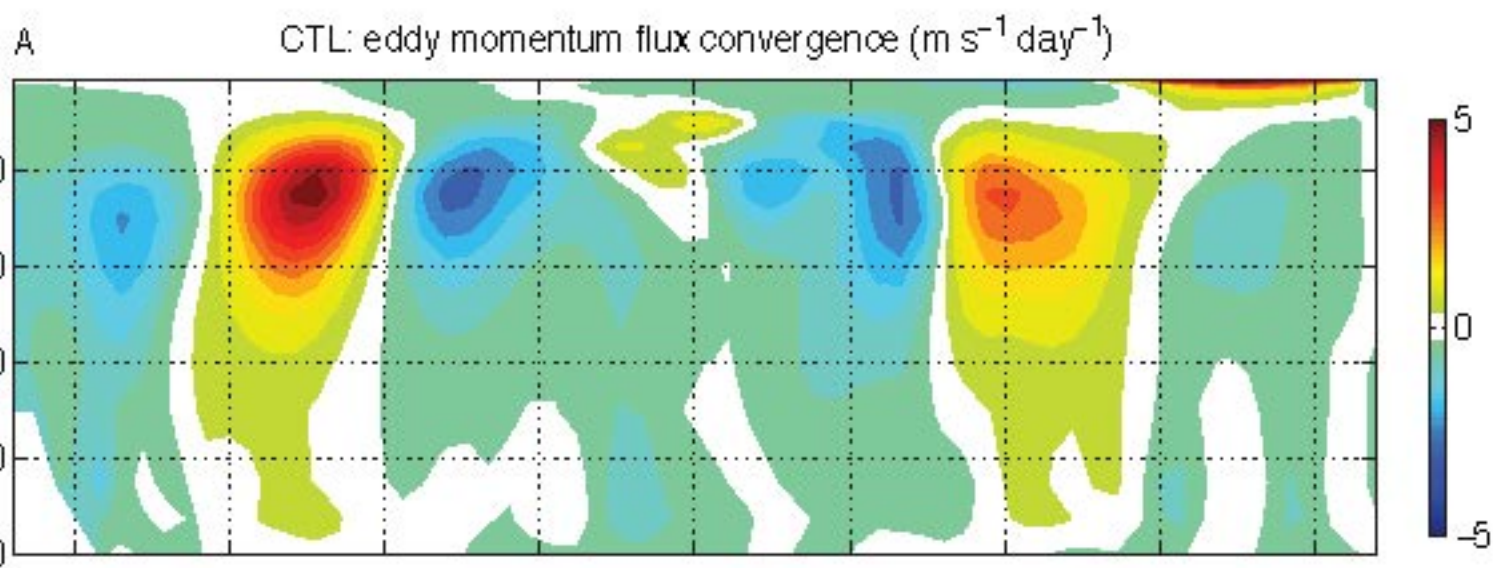
G: RAD



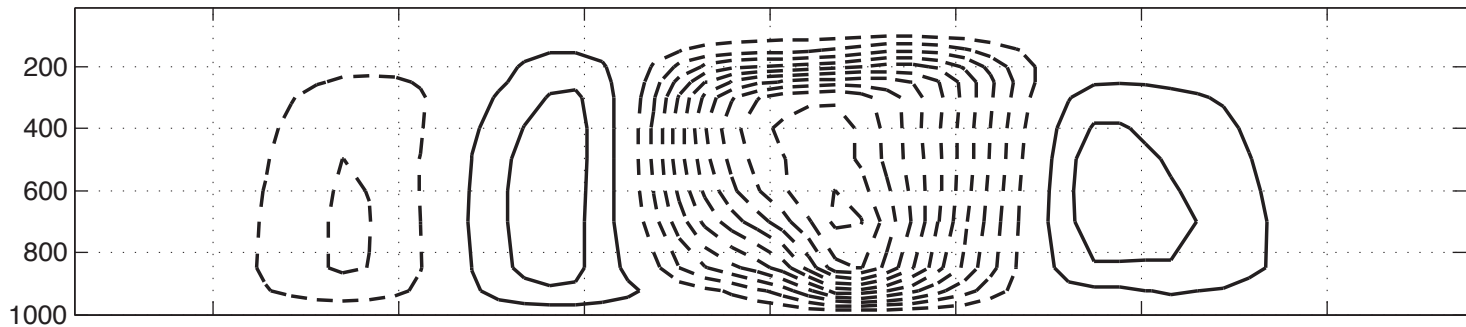
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zonal wavenumber

-10 -5 0 5 10
zonal wavenumber

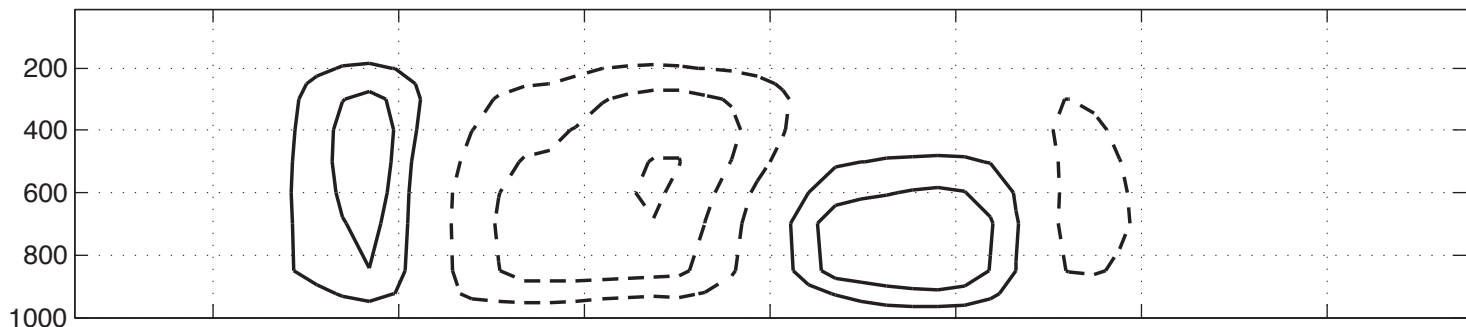
-10 -5 0 5 10
zonal wavenumber



A CTL: Mass flux streamfunction (interval: $2 \times 10^{10} \text{ kg s}^{-1}$)



B CHNn: Mass flux streamfunction anomaly (interval: $2 \times 10^{10} \text{ kg s}^{-1}$)



C CHN: Mass flux streamfunction anomaly (interval: $2 \times 10^{10} \text{ kg s}^{-1}$)

